


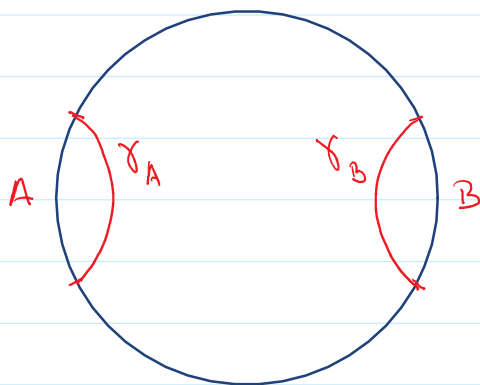
We can now compute holographically The mutual information for a bipartite system  $\mathcal{H}_A \otimes \mathcal{H}_B$ :

$$I(A:B) = S_A + S_B - S_{AB}$$

This is positive because of subadditivity of entanglement entropy,  $S_A + S_B \geq S_{AB}$  

[which in QFT, for  $A \cap B \neq \emptyset$  is satisfied because  
The UV divergences on the lhs are at least as  
big as the ones in the rhs]

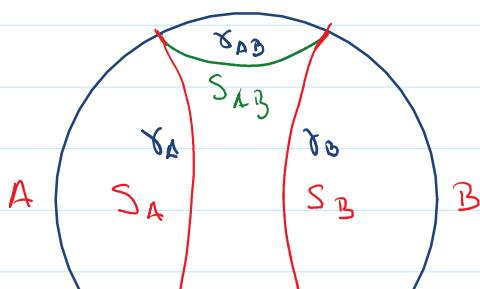
Consider the mutual info between two regions  $A, B$  in a holographic set up



$$S_{AB} = S_A + S_B$$

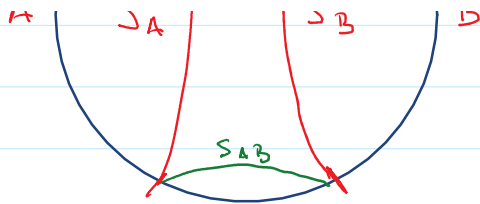
$$\Rightarrow I(A:B) = 0$$

but when  $A$  and  $B$  become large enough, the RT surface for  $A \cup B$  changes:



$$I(A:B) = S_A + S_B - S_{AB} \neq 0$$

This means there are correlations between  $A$  and  $B$ , and the



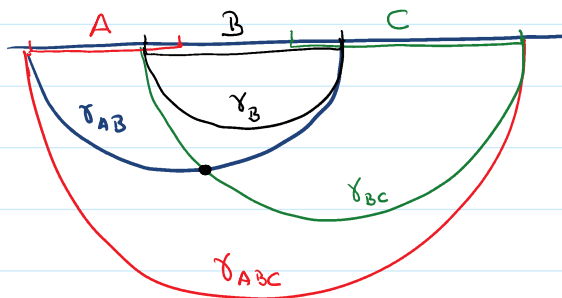
$$\text{Area}(\gamma_A) + \text{Area}(\gamma_B) > \text{Area}(\gamma_{AB})$$

between A and B, and the combined system contains more information than A and B separately  
i.e. more information about how to reconstruct the bulk

Other properties of entanglement entropy become nicely geometrized by the HEE

Eg Strong subadditivity for a tripartite system AUBUC

$$S_{AB} + S_{BC} - S_B \geq S_{ABC}$$



At the point where  $\gamma_{AB}$  and  $\gamma_{BC}$  intersect, split and rejoin and then deform to  $\gamma_{ABC}$  and  $\gamma_B$ .

Then, since  $\gamma_{ABC}$  and  $\gamma_B$  are minimal, it follows that

$$\text{Area}(\gamma_{AB}) + \text{Area}(\gamma_{BC}) \geq \text{Area}(\gamma_{ABC}) + \text{Area}(\gamma_B)$$

## - Entanglement wedge

QFT state in a region A is described by a density matrix  $\rho_A$ .

Q: What is the region of the bulk that can be reconstructed from knowledge of  $\rho_A$ ?

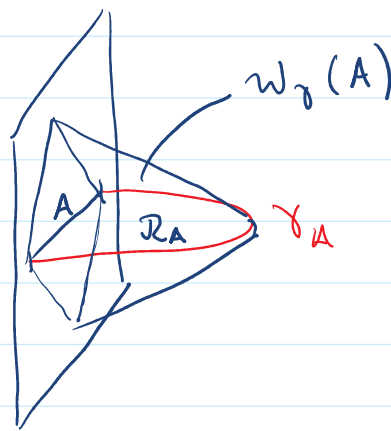
A: The entanglement wedge  $W_\gamma(A)$

Given  $\rho_A$ , it is possible to compute  $S(A)$ .

Holographically, this is given by the area of the extremal surface  $\gamma_A$ .

The homology region  $R_A$  is enclosed between  $\gamma_A$  and  $A$ . The entanglement wedge  $W_\gamma(A)$  is the bulk domain of dependence of  $R_A$ :

$$W_\gamma(A) = D(R_A)$$

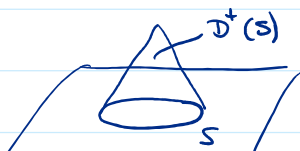


Reminders in GR: for an achronal set  $S$  (ie no two points in  $S$  are timelike-connected) we define:

The future domain of dependence of  $S$ :

$$D^+(S) = \{ p \in M \mid \text{every past inextendible causal curve through } p \text{ intersects } S \}$$

$\downarrow$  not ending       $\downarrow$  Timelike or null

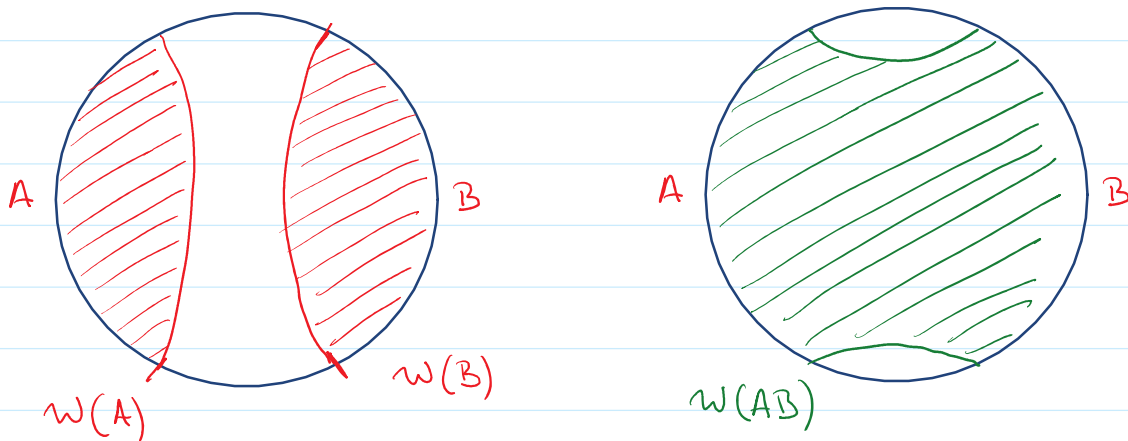




Similarly define past domain of dependence  $D^-(S)$   
and then domain of dependence  $D(S) = D^+(S) \cup D^-(S)$

"set of events that can be determined from data on S"

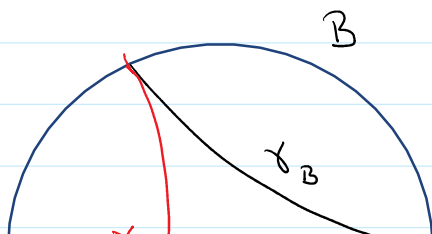
We can see that when  $I(A:B) > 0$ , one can reconstruct a larger region of the bulk with  $A \cup B$  than with  $A$  and  $B$  separately:



### Holographic quantum error correction

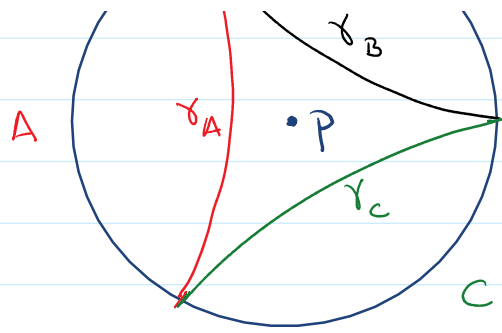
$\partial \text{AdS} = A \cup B \cup C$        $A, B, C$  non overlapping

$\exists$  point  $p$  in the bulk cannot be reconstructed from either of  $A, B, C$  alone, but it can be from the union of any pair of them:



$p \notin W(A), W(B), W(C)$

$p \in W(AB), W(AC), W(BC)$



$$p \in w(AB), w(AC), w(BC)$$

We could erase either of A, B, C and still reconstruct p

"Spacetime is a holographic quantum error-correcting code"